



# DARK MATTER

## APPLIED IN STRANGE STARS

LASNPA XIV/2024

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# Introduction

- Our work is focused on strongly interacting particles that are studied by Quantum Chromodynamics (QCD);
- In Quantum Chromodynamics we have quarks and gluons that are studied under different conditions of temperature and density (or through the chemical potential);
- There are two possible states of matter: the hadronic one and the quark and gluon plasma (QGP)

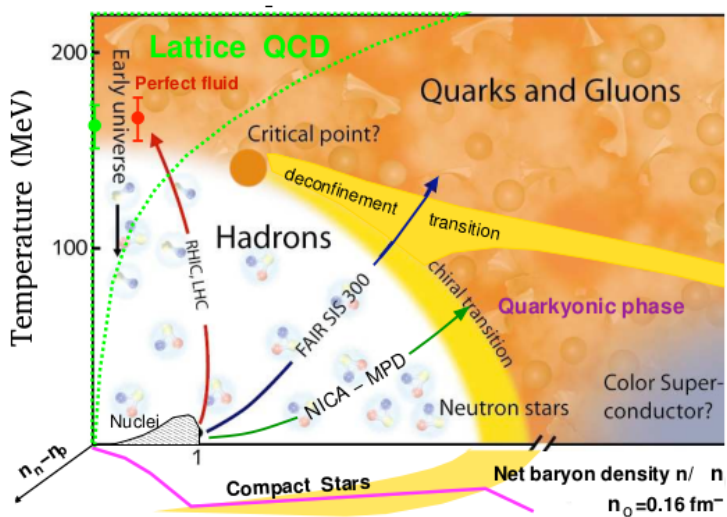


Figure: QCD phase diagram.

- The difficulty of studying QCD at the non-perturbative regime requests different approaches:
  - Lattice QCD (LQCD);
  - Dyson-Schwinger equations;
  - Effective/ phenomenological quark/gluons models:
    - Massachusetts Institute of Technology (MIT) bag model;
    - Nambu-Jona-Lasino (NJL) model;
    - The Equiparticle (EQP) model;

# The equiparticle (EQP) model

- Let's begin by understanding how the Equiparticle model works;
- In the 70's, the Bodmer-Witten Hypothesis was proposed: the absolutely stable fundamental state of matter would be strange quark matter (SQM);
- The EQP model was proposed in 2014 in order to study SQM by being a density and/or temperature dependent model;
- This model corrects the lack of thermodynamic consistency present on other models of the same type.

The relevant equations describing the model are given by:

$$m_i = m_{i0} + m_l = m_{i0} + \frac{D}{\rho_b^{1/3}} + C\rho_b^{1/3}$$

$$p_{SQM} = -\Omega_0 + \rho_b \frac{\partial m_l}{\partial \rho_b} \frac{\partial \Omega_0}{\partial m_l} \quad \epsilon_{SQM} = \Omega_0 - \sum_i \mu_i^* \frac{\partial \Omega_0}{\partial \mu_i^*}$$

$$\Omega_0 = - \sum_i \frac{\gamma}{24\pi^2} \left[ \mu_i^* k_{Fi} \left( k_{Fi}^2 - \frac{3}{2} m_i^2 \right) + \frac{3}{2} m_i^4 \ln \frac{\mu_i^* + k_{Fi}}{m_i} \right]$$

$$\mu_i = \mu_i^* + \frac{1}{3} \frac{\partial m_l}{\partial \rho_b} \frac{\partial \Omega_0}{\partial m_l} \quad k_{Fi} = \sqrt{\mu_i^{*2} - m_i^2} \quad \rho_i = \frac{\gamma k_{Fi}^3}{6\pi^2} \quad \rho_b = \frac{1}{3} \sum_i \rho_i$$

$$\Omega_{SQM} = -p_{SQM}$$

where  $i = u, d, s$  and  $\gamma = 6$ ;

# Thermodynamic consistency

The conditions that guarantee the thermodynamic consistency of the model are:

$$\Delta_i \equiv \left. \frac{dS}{d\rho_i} \right|_{T, \{\rho_{k \neq i}\}} - \left. \frac{d\mu_i}{dT} \right|_{T, \{\rho_k\}} = 0$$

$$\Delta_{ij} \equiv \left. \frac{d\mu_i}{d\rho_j} \right|_{T, \{\rho_{k \neq j}\}} - \left. \frac{d\mu_j}{d\rho_i} \right|_{T, \{\rho_{k \neq i}\}} = 0$$

$$\Delta = p - \rho^2 \left. \frac{d}{d\rho} \left( \frac{F}{\rho} \right) \right|_T = 0$$



# Stellar matter

- In order to describe strange stars we have to consider some features of stellar matter:
  - Zero temperature and high density;
  - Beta equilibrium:  $\mu_u + \mu_e = \mu_d = \mu_s$  and  $\mu_u^* + \mu_e = \mu_d^* = \mu_s^*$ ;
  - Weak interactions:  $d, s \leftrightarrow u + e + \nu_e$ ;
  - Charge neutrality:  $\frac{2}{3}\rho_u - \frac{1}{3}\rho_d - \frac{1}{3}\rho_s - \rho_e = 0$ ;
  - EOS with electrons contribution:

$$\varepsilon = \varepsilon_{SQM} + \frac{\mu_e^4}{4\pi^2}, \quad \text{and} \quad p = p_{SQM} + \frac{\mu_e^4}{12\pi^2},$$

- with  $\mu_e = (3\pi^2\rho_e)^{1/3}$ ;
- $p_{SQM}$  and  $\varepsilon_{SQM}$  are the EOS of quarks contribution.

# Stability window

- According to Bodmer-Witten hypothesis, in order to SQM to be absolutely stable, it needs to fulfill the following conditions:
  - $(E/A)_{\min} = (\mathcal{E}/\rho_b)_{\min} \leq 930$  MeV for SQM;
  - $(E/A)_{\min} = (\mathcal{E}/\rho_b)_{\min} > 930$  MeV for 2QM;
- Quarks current mass are the ones most recent data provided by the Particle Data Group (PDG):
  - For stellar matter:  $m_{u0} = 2.16$  MeV,  $m_{d0} = 5.15$  MeV, and  $m_{s0} = 90$  MeV.
- Applying all the above we constructed the stability window for stellar matter.

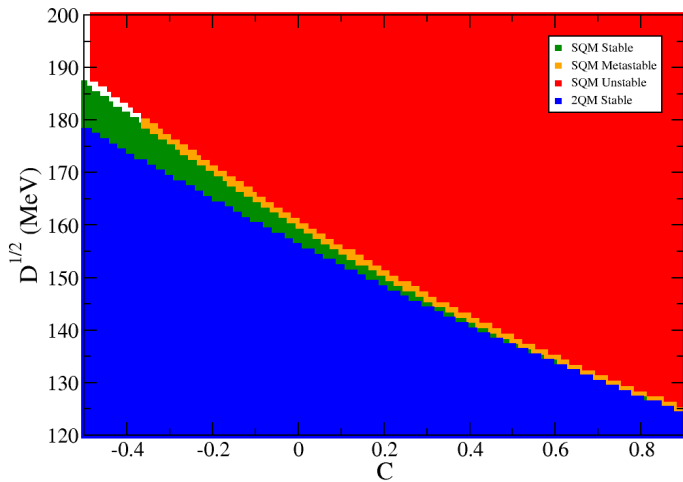


Figure: Stability window for stellar matter: EQP model.

# Dark Matter

- Dark matter (DM) is thought to permeate throughout the universe, exerting gravitational effects on visible matter;
- It could be possible to have DM particles trapped within the intense gravitational fields of compact stars due to:
  - Its capture during formation;
  - Accretion;
  - Annihilation and heating;
- As a consequence, the presence of DM inside compact stars could influence its properties;
- In our current work we introduce the dark matter as a bosonic fluid in the EQP model in order to study its presence in strange stars;

# Bosonic asymmetric dark matter

- Baryogenesis period: during the creation of matter-antimatter a slight asymmetry in the production of matter over antimatter occurred (allowing the formation of all the visible matter in the universe, such as: stars, galaxies, etc);
- Asymmetric dark matter could have occurred from a similar process (with DM being slight higher than anti-DM due to some fundamental differences in the interactions of dark matter particles compared to ordinary matter particles);
- This excess dark matter survived annihilation with dark antimatter and became the dark matter that we study today;
- This feature explains why DM interacts differently with ordinary matter;

- Compact stars are modeled from its equations of state introduced in the Tolman-Oppenheimer-Volkoff (TOV) equations;
- TOV equations provide us the mass-radius relation of compact stars that is calculated through the central density;
- To calculate the structure of compact stars with asymmetric dark matter, we treat the system as being composed by two separate fluids that only interact through gravity;

The equations concerning the two-fluid formalism are:

$$p(r) = p_{SQM}(r) + p_{\chi}(r) \quad \epsilon(r) = \epsilon_{SQM}(r) + \epsilon_{\chi}(r)$$

$$\frac{d\alpha(r)}{dr} = \frac{1}{c^2} \frac{Gc^2M(r) + 4\pi r^3 G\rho(r)}{r[rc^2 - 2GM(r)]}$$

$$\frac{dp_{SQM}}{dr} = -(\epsilon_{SQM} + p_{SQM}) \frac{d\alpha(r)}{dr} \quad \frac{dp_{\chi}}{dr} = -(\epsilon_{\chi} + p_{\chi}) \frac{d\alpha(r)}{dr}$$

$$\frac{dM_{SQM}(r)}{dr} = 4\pi r^2 \frac{\epsilon_{SQM}(r)}{c^2} \quad \frac{dM_{\chi}(r)}{dr} = 4\pi r^2 \frac{\epsilon_{\chi}(r)}{c^2}$$

where  $M_{\chi}(r)$  = gravitational mass of ADM,  $M_{SQM}(r)$  = gravitational mass of baryonic matter (SQM in this case) and  $M(r) = M_{SQM}(r) + M_{\chi}(r)$ ;

We also need to define the ADM mass-fraction,  $F_\chi$ , which is, the relative amount of gravitational mass of ADM compared to the total gravitational mass of the compact star:

$$F_\chi = \frac{M_\chi(R_\chi)}{M_\chi(R_\chi) + M_{SQM}(R_{SQM})}$$

where  $M_\chi(R_\chi)$  = the total accumulated ADM gravitational mass evaluated at the ADM core radius and  $M_{SQM}(R_{SQM})$  = the baryonic matter gravitational mass evaluated at the baryonic radius.



Following, the EOS for the bosonic ADM are constructed to be:

$$\epsilon_{\chi}(r) = m_{\chi}c^2n_{\chi} + \frac{g_{\chi}^2}{2m_{\phi}^2} \frac{\hbar^3}{c} n_{\chi}^2,$$

$$p_{\chi}(r) = -\epsilon_{\chi} + \tilde{\mu}_{\chi}n_{\chi} = \frac{g_{\chi}^2}{2m_{\phi}^2} \frac{\hbar^3}{c} n_{\chi}^2,$$

$$\tilde{\mu}_{\chi}(r) = \frac{\partial \epsilon_{\chi}}{\partial n_{\chi}}$$

where  $m_{\chi}$  = ADM particle mass,  $n_{\chi}$  = number density of ADM,  $\tilde{\mu}_{\chi}$  = ADM chemical potential (local Lorentz frame), and  $c$ ,  $\hbar$  and  $G$  are treated in natural units.

# Constraints of ADM

The bosonic ADM EOS constraints are:

- To avoid the ADM particles from escaping the compact star  $m_\chi \geq 10^{-2}$  MeV and  $m_\chi \leq 10^8$  MeV (we chose  $m_\chi = 15 \cdot 10^3$  MeV);
- The effective self-interaction strength are captured in the range  $10^{-2} \leq \frac{g_\chi}{m_\phi/\text{MeV}} \leq 10^3$  (we chose  $\frac{g_\chi}{m_\phi/\text{MeV}} = 0.1$ );
- In order to obtain the EOS of the dark matter contribution we consider  $\epsilon_\chi(r) = f \cdot \epsilon_{SQM}$ , where  $f$  indicates the fraction of dark matter energy density present in the strange quark matter energy density ( $f$  varies from 0.1 to 1.0);

# Results

- The observational data in the mass-radius profiles are the ones provided by LIGO and Virgo Collaboration, and the NICER mission concerning the millisecond pulsars PSR J0030+0451, and PSR J0740+6620;
- We also include the recent observational data concerning the measurement of HESS J1731-347;

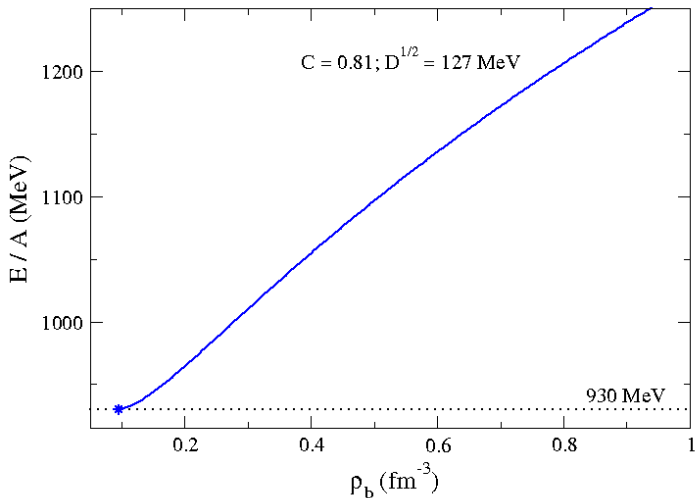


Figure: Thermodynamic consistency for stellar matter: EQP model.

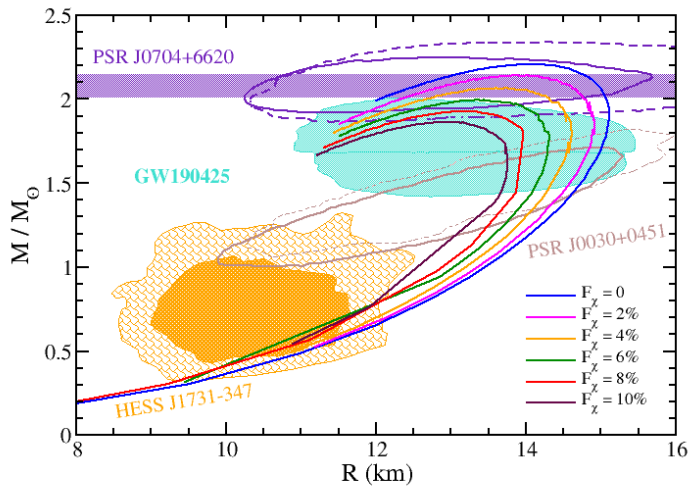


Figure: Mass-Radius profiles.

# Future perspectives

- Apply the fermionic dark matter in the EQP model as proposed in the recent article "Exploring robust correlations between fermionic dark matter model parameters and neutron star properties: A two-fluid perspective";
- Discuss the results obtained from both proposals;

# References

List of references



Last published article



