Constraints on magnetars as the engines of prompt emission in GRBs

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Which central engine?

Both long and short GRBs may involve either a BH or NS central engine (CE) Traditionally, energetics considered as clearest imprint of CE Coupling the CE with dissipation models is crucial for going further

Rapidly rotating magnetar



Black hole



PB & Giannios 2017 Gradual magnetic dissipation

MHD outflow arranged in a striped wind configuration

 $\lambda = \frac{\pi c}{\Omega} \approx 10^8$ cm for magnetars







Magnetic energy dissipated by reconnection and is converted to bulk kinetic energy and acceleration of relativistic particles

 $\Gamma = \sigma_0 \left(\frac{r}{r_0}\right)$

Drenkhahn 02

Energy per baryon

Saturation radius= $\lambda \sigma_0^2/6\epsilon$

Particle acceleration in high magnetization

Particle spectrum becomes harder with increasing magnetization



Sironi & Spitkovski 14

Gradual magnetic dissipation

Dissipated energy converted to radiation

Sub-Photospheric dissipation reprocessed to quasi thermal emission

Optically thin synchrotron

Thermal emist

Super-Photospheric dissipation dominated by synchrotron

PB & Giannios 2017

Results

Four model parameters:

L (jet luminosity), $\frac{\lambda}{\varepsilon}$ (scale of striped wind over reconnection outflow velocity), σ_0 (energy per baryon), ξ fraction of accelerated electrons



Predictions / Observables

Efficiency of 0.1-0.2 consistent with afterglow observations (PB, Nava, Barniol Duran, Piran 15)

Natural preferred range of Lorentz factors

 $290 \left(\frac{\lambda}{\varepsilon}\right)_{8}^{-1/5} L_{52}^{1/5} \lesssim \Gamma \lesssim 1800 \left(\frac{\lambda}{\varepsilon}\right)_{8}^{-1/5} L_{52}^{1/5} \left(\frac{\varepsilon_{e}}{\xi}\right)^{1/5}$ For lower values

 $r_{ph} > r_s$ emission completely thermalized + adiabatic losses For larger values electrons are slow cooling and inefficiency is large

Bursts with softer α have weaker thermal bump and observable to higher frequencies

Emission in X-ray and optical self absorbed and consistent with upper limits from observations (PB & Piran 14)

Material ejected with lower L, σ_0 would produce X-ray flares that are self-absorbed in the optical band (PB & Kumar 16)

Sensitivity to initial magnetization

The observed luminosity and peak energy strongly depend on the energy per baryon, σ_0 and its evolution $\nu_c(r_s) \propto \sigma_0^9$!!! While $\nu_m(r_s) \propto \sigma_0^{-2}$



For $\sigma_0 \ge 2 \times 10^3$ energy dissipation occurs in strongly slow cooling conditions and resulting GRB is very faint

Remains true for general models with $\Gamma \propto r^n$

Furthermore when n<1/4 or n>2/3 E_p and the luminosity evolve strongly with σ_0 (and therefore also with time), contrary to observations



Magnetar models are more strongly constrained *(for better or worse)*

- $\lambda = \frac{\pi c}{\Omega} \approx 10^8 \text{cm}$ for magnetars
 - 1. Energy Strong energy limits $\sim 10^{51}$ erg from stational energy reservoir
 - 2. Length scale Evolution of spin frequency due to dipole radiation – natural scale λ for flipping of magnetic field

3. Energy per baryon - related to neutrino driven mass loss (Metzger et al. 10)

For BHs $\lambda = 10 - 10^3 r_g \approx 10^7 - 10^9$ cm

- 1. Energy basically limitless. 10^{51} erg requires accreting only $0.01M_{\odot}$ at 10% efficiency of $\dot{M}c^2$
- 2. Length scale Evolution of λ unknown (could be chaotic and likely determined by the disc)
- Energy per baryon No robust prediction for the energy per baryon (yet)

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Magnetar Wind Model

Energy loss dictated by dipole spin-down

Mass loss driven by neutrinos from cooling NS atmosphere ~ 1 min At later times NS transparent to ν and mass loss given by pair creation from the vacuum electric fields

Wind characterized by magnetization $\sigma_0 = \frac{\dot{E}}{\dot{M}c^2} \ge \Gamma$

Before SN Shock Launch

After Shock Launch

vetn-pte

Neutrino-Heated Wind

Burrows, Hayes, & Fryxell 1995

Magnetar Wind Model Transition to vacuum pairs –

Energy per electrons decreases by ~1000!



Magnetar Wind Model



• For $B \approx 10^{16} G$ most of the energy released when $20 < \sigma_0 < 3000$

• Available energy for GRB $\lesssim 0.25 E_{rot}$, with a maximum at $3 \cdot 10^{51}$ erg

Coupling to prompt emission models

Dissipation models required to constrain prompt emission

Consider three generic possibilities:
Pure fireball – Dissipation close to central engine
Internal shocks – Dissipation at large distances
Gradual Magnetic dissipation – Dissipation across wide range of radii (PB & Giannios 2017)



Pure Fireball

Strong temporal evolution, hard spectrum, shallow late time decay



Internal Shocks Strong temporal evolution, soft spectrum, inefficient



Gradual magnetic dissipation Stable E_p, L_p, reasonable spectrum, very steep late time decay



Duration distribution

Unless dissipation completed below the photosphere, $t_{GRB} = Min(t_{SD}, t_{\sigma_0}) \leq 100sec$ Short GRB if $t_{collapse} < t_{SD}$ - similar luminosity, smaller energy – remnant is BH Ultra long GRBs implausible from this engine (Fallback accretion? Work in progress)



Energy distribution

Limited by energy released from $\max[t_b, t(\sigma = 100)]$ and t_{σ_0} as well as efficiency of prompt GRB mechanism Typical energies easy to produce, but some GRBs have 30 times larger than obtainable from magnetars



Data from Goldstein et al. 2016



Magnetar model provides robust predictions

- Magnetization = σ_0 is a crucial parameter
- GRB available energy $\leq 0.25 E_{rot}$, with a maximum at $3 \cdot 10^{51}$ erg
- Fireball and Internal shocks coupled with magnetar engines in contention with observations
- Gradual dissipation model provides realistic energies, time-scales, spectra, temporal evolution, etc. with no fine tuning but still cannot easily account for very long or very energetic GRBs

Thank you!

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